

# Comment on "Aging Effects in a Lennard-Jones Glass"

In a recent Letter Kob and Barrat<sup>1</sup> reported results of molecular dynamics simulations for the off-equilibrium dynamics in a binary Lennard-Jones (LJ) glass. The main conclusions of their work was 1) they find aging in this glassy systems and 2) that they find a *simple* aging scenario close to a  $t/t_w$  scaling, which is very reminiscent of comparable studies in spin glasses. In this comment we would like to emphasize that a different aging scenario, known under the name *activated dynamics* scaling, is much more appropriate for the system under consideration than the one proposed by Kob and Barrat<sup>1</sup>.

For this reason we repeated the simulation by Kob and Barrat, using exactly the same potential (Lennard-Jones for a binary mixture), the same parameters (same diameters, mixture, density and temperatures) and the same quenching procedure ( $T_i=5$ ,  $T_f=0.4$ ) however with much larger systems ( $32768 = 32^3$  particles) and similar times ( $2 \cdot 10^6$  time steps, 1 time step corresponding to 0.01 LJ-units). The aging properties of the system manifest themselves in the two-time autocorrelation function

$$C_q(t + t_w, t_w) = \frac{1}{N} \sum_i e^{i \cdot q \cdot [r_i(t+t_w) - r_i(t_w)]}, \quad (1)$$

where  $r_i(t)$  is the position of particle  $i$  at time  $t$  and the absolute value of  $q$  corresponds to the first maximum in the structure function. We choose 100 randomly distributed vectors and averaged  $C_q$  over these vectors. The function (1) was evaluated after every 10 time steps and  $5^n$  measurements were averaged over to improve statistics. We convinced ourselves that different quenching procedures with identical initial and final temperatures,  $T_i$  and  $T_f$ , lead to the same scaling behavior.

In [1] it has been suggested that  $C_q(t + t_w, t_w)$  obeys

$$C_q(t + t_w, t_w) \sim \tilde{C}(t/t_r) \quad (2)$$

with a relaxation time  $t_r \propto t_w^\alpha$ . We checked this *Ansatz* for our data and display the result in the inset of Fig. 1, surprisingly we find an exponent  $\alpha \sim 1.1$ , very close to one (corresponding to simple  $t/t_w$  scaling) but different from the one  $\alpha = 0.88$  reported in [1]. The data collapse in the asymptotic regime is not at all satisfying, the data for different waiting times coincide exactly only for  $C_q = 0.45$ . For this reason we tried another aging scenario, proposed in the context of spin glasses by Fisher and Huse<sup>2</sup>, which we call the *activated dynamics*:

$$C_q(t + t_w, t_w) \sim \tilde{C}\{\ln((t + t_w)/\tau)/\ln(t_w/\tau)\} \quad (3)$$

where  $\tau$  is a fit-parameter and plays the role of an effective microscopic time scale. Fig.1 we show the scaling plot for such a scenario, which gives a much better data collapse in the asymptotic regime  $t \geq t_w$ .

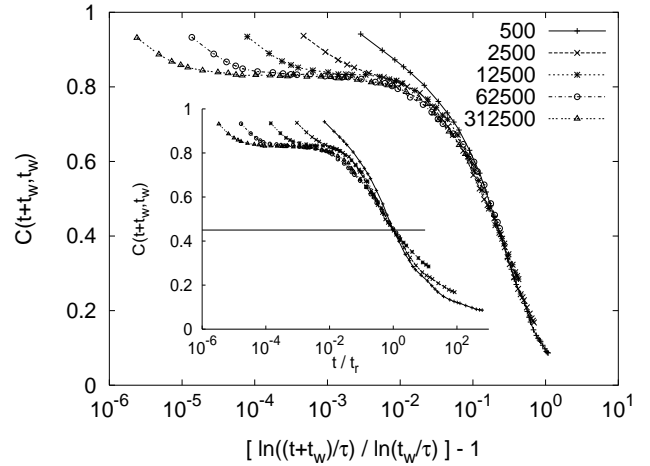


FIG. 1. Activated dynamics scaling plot according to eq.(3) with  $\tau = 0.5$ . Note that we shifted the scaling variable by 1 to the left to have a better resolution of the crossover region. The inset shows a scaling plot of our data according to the scenario proposed by Kob and Barrat [1], see eq.(2), the full line corresponds to  $C_q = 0.45$ . The relaxation times are  $t_r=1450$ , 10600, 70000, 600000 and 3000000 for  $t_w=500$ , 2500, 12500, 62500 and 312500, roughly a dependence  $t_r \propto t_w^{1.1}$ .

The origin of such an activated dynamics scaling in spin glass phenomenology<sup>2</sup> is simply a logarithmically slow coarsening process  $\xi(t) \sim \ln(t)^a$ , where  $\xi(t)$  is a time dependent spatial correlation length and  $a$  some exponent. This plus the observation that in coarsening dynamics the two time correlation function  $C_q(t + t_w, t_w)$  should depend on the ration of the two length scale  $\xi(t_w)/\xi(t + t_w)$  alone yields the aging behavior (3).

Three things are worth being noted: 1) In the context, in which (3) was first suggested, namely the 3d EA spin glass, this form does not seem to work<sup>3</sup>. 2) Only very recently a growing length scale has been observed in the very same model we are considering here<sup>4</sup>. 3) An even better data collapse can be obtained by plotting  $C_q(t + t_w, t_w)$  versus  $\ln(t)/\ln(t_r)$ , with a relaxation time  $t_r$  individually chosen for each waiting time  $t_w$ . Here it turns out that  $t_r(t_w)$  grows faster than with a power law.

To conclude we have shown that the aging behavior of a Lennard-Jones glass is more appropriately described by an activated dynamics scaling rather than simple aging, as claimed by Kob and Barrat in [1].

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<sup>2</sup> D. S. Fisher and D. A. Huse, Phys. Rev. B **38** (1988) 386 and Phys. Rev. B **38** (1988) 373.

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<sup>4</sup> D. Lancaster and G. Parisi, J. Phys. A **30** (1997) 5911; G. Parisi, cond-mat/9801034.